Optimal Control in Transition Path Theory

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Amar Shah (University of California, Berkele Optimal Control in Transition Path Theory



- 2 Optimal Control
- Finding an Optimal Controller
- 4 Numerical Results

Introduction

- 2 Optimal Control
- 3 Finding an Optimal Controller
- 4 Numerical Results

Ordinary Differential Equation

$$dx = \mathbf{b}(x)dt$$

(1)

We can solve this differential equation to understand how our system evolves in time. But what if we added randomness to our system:

Stochastic Differential Equation

$$dx = \mathbf{b}(x)dt + \sigma(\mathbf{x})d\mathbf{w}$$

(2)

Goal

We have two compact, connected sets $A, B \subseteq \mathbb{R}^n$ and we want to consider the paths in our evolution that start in A and end up in B.



Figure: E & Vanden-Eignden(2006, 2010) [1] [2]

We call these reactive trajectories.

Forward Committor

We define the forward committor $q^+(x)$ as:

$$\tau_A(x) = \min\{t > 0 : x(0) = x; x(t) \in \delta A\}$$

$$\tau_B(x) = \min\{t > 0 : x(0) = x; x(t) \in \delta B\}$$

$$q^+(x) = \mathcal{P}_x[\tau_B(x) < \tau_A(x)]$$

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3 Finding an Optimal Controller



In a dynamic system where there is some function v that "we" can control. Our goal is to pick $v(\cdot)$ to maximize the pay off function $P[v(\cdot)]$



Figure: Evans(1980) [3]

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Let's return to our setting at the beginning of the presentation:

Stochastic Differential Equation

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Stochastic Differential Equation

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We want to introduce a control function $v : \mathbb{R}^n \to \mathbb{R}$ and we have the equation:

Controlled Stochastic Differential Equation

 $dx = (\mathbf{b}(x) + \mathbf{v}(x))dt + \sigma d\mathbf{w}$





Finding an Optimal Controller

4 Numerical Results

Gao et al Theorem 3.3[4]: If we have the controlled SDE $dx = (-\nabla U + \mathbf{v}(x))dt + \sqrt{2\epsilon}d\mathbf{w}$, we have the optimal control $v^* = \frac{2\epsilon\nabla q}{a}$

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But remember our original SDE:

Controlled SDE

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How can we apply Gao's argument to this general case?

Maybe we can't apply Gao's Argument to the full general case, but we can definitely do it for general drift ${\bf b}$ and invertible σ :

Controlled SDE

$$dx = (\mathbf{b}(x) + \mathbf{v}(x))dt + \sigma d\mathbf{w}$$

To solve it for the general case, we need to solve the Hamilton Jacobi Bellman equation.

We introduce the controlled SDE, with control v(x):

$$\begin{cases} dX(s) = f(X(s), v(X_s))dt + \sigma(X_s)dW(s) \ (t \le s \le T) \\ X(t) = x \end{cases}$$

We have the expected payoff function:
 $P_{x,t}[v(\cdot)] := E[\int_t^T r(X(s), v(X_s))ds + g(X(T))]$
The value function is $\gamma(x, t) = \sup_{v(\cdot) \in \mathcal{V}} P_{x,t}[v(\cdot)]$, i.e. the best possible

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Stochastic-Hamilton-Jacobi-Bellman equation

Then our value function γ solves this PDE: $\begin{cases} \gamma_t(x,t) + \frac{\sigma^2}{2} \Delta v(x,t) + \max_{v \in \mathcal{V}} \{f(x,v) \cdot \nabla_x \gamma(x,t) + r(x,a)\} = 0 \\ \gamma(x,T) = g(x) \end{cases}$

Solution

Theorem

The optimal controller for the general SDE when σ is invertible is $\mathbf{v}^* = \sigma \sigma^T \frac{\nabla q}{q}$

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• Solve the maximization term to get $\mathbf{v}=-\nabla\gamma$

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- Solve the maximization term to get $\mathbf{v}=-\nabla\gamma$
- Plug in for v to simplify our HJB to:

$$\frac{1}{2}\sigma\sigma^{\mathsf{T}}:\nabla\nabla\gamma+b\cdot\nabla\gamma-\frac{1}{2}(\nabla\gamma)^{\mathsf{T}}\sigma\sigma^{\mathsf{T}}\nabla\gamma=0$$

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• Show that this is satisifed by $\gamma = -log(q)$ using properties from Transition Path Theory

Full Langevin is given by the equation:

$$\begin{cases} d\mathbf{x} = m^{-1}\mathbf{p}dt \\ d\mathbf{p} = -(\nabla U + \gamma \mathbf{p})dt + \sqrt{2\gamma\beta^{-1}m}d\mathbf{w} \end{cases}$$
(3)
Thus, $\sigma = \sqrt{2\gamma\beta^{-1}m} \begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix}$ which is not invertible.

Theorem

The optimal controller for the general SDE when σ is (symmetric) and invertible is $v^* = 2\gamma\beta^{-1}m\frac{\nabla_p q}{q}$ applied to the momentum part of the equation.

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Overdamped Langevin dynamics with Mueller Potential



Figure: 10 Uncontrolled and Controlled trajectories until T = 10

Overdamped Langevin dynamics with Mueller Potential



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This follows the overdamped equation $dx = (-\nabla U(x) + \frac{2}{\beta} \frac{\nabla q}{q})dt + \sqrt{\frac{2}{\beta}} dw$. In the uncontrolled case, we have a success rate of 0 and in the controlled case we have a success rate of 0.932



Figure: 25 Uncontrolled and Controlled trajectories until T = 10 for $\beta = 10$

- The Duffing Oscillator is essentially Full Langevin Dynamics in two dimensions.
- The uncontrolled process transitions with rate 0.313 and the controlled process transitions with rate 0.917

Results

- **2** Proved an Optimal Controller result for Full Langevin dynamics, where σ is not invertible
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Further Work

- Apply these results to more complicated systems including Full Langevin Dynamics in multiple dimensions
- Apply these ideas to similar problems, for example to calculate escape rates.



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